Many Electron Spin Eigenfunctions

An arbitrary Slater determinant for N electrons can be written as

$$\hat{\mathcal{A}}\Phi(1,2,\cdots,N)\chi_{M}(1,2,\cdots,N)$$

Where $\Phi(1,2,\dots,N) = a(1)b(2)\dots c(N)$ is a product of N orthonormal spatial functions and $\chi_M(1,2,\dots,N)$ is a product of $N_\alpha \alpha$ spin functions and $N_\beta \beta$ spin functions and therefore has an eigenvalue of \hat{S}_z equal to $M = \frac{1}{2}(N_\alpha - N_\beta)$. Note that we do not consider double occupied spatial orbital when investigating the spin characteristics of a Slater Determinant. To determine the effect of \hat{S}^2 on the Slater determinant we need only investigate $\hat{S}^2 \chi_M(1,2,\dots,N)$. Accordingly if we want a Slater determinant to be an eigenfunction of \hat{S}^2 with an eigenvalue S(S+1) we need only to write

 $\hat{\mathcal{A}}\Phi(1,2,\cdots,N)\chi_{SM}(1,2,\cdots,N)$

Where $\hat{S}^2 \chi_{SM}(1,2,\dots,N) = S(S+1)\chi_{SM}(1,2,\dots,N)$. The problem is how to find $\chi_{SM}(1,2,\dots,N)$ and there are several general approaches. Consider for example a three- electron system. We may form $2^3 = 8$ spin functions and we list them below labeled by the M quantum number. We assume the electron order is 1, 2, 3, so $\alpha\beta\alpha = \alpha(1)\beta(2)\alpha(3)$

Number	Function	М
1	ααα	3/2
2	ααβ	1/2
3	αβα	1/2
4	βαα	1/2
5	ββα	-1/2
6	βαβ	-1/2
7	lphaetaeta	-1/2
8	βββ	-3/2

Note that $\alpha\alpha\alpha & \beta\beta\beta$ must be eigenfunctions of \hat{S}^2 with S=3/2 and $M = \pm 3/2$. Functions 2, 3 and 4 must span the space containing one quartet (S=3/2) and two doublets (S=1/2) all with M=1/2 while functions 5, 6, and 7 lie in the M = -1/2 subspace of the eigenfunctions of \hat{S}^2 . So for the M=1/2 subspace we may write

$$\varphi_{1} = |S = 3/2, M = 1/2\rangle = a_{1}\alpha\alpha\beta + a_{2}\alpha\beta\alpha + a_{3}\beta\alpha\alpha$$
$$\varphi_{2} = |S = 1/2, M = 1/2\rangle = b_{1}\alpha\alpha\beta + b_{2}\alpha\beta\alpha + b_{3}\beta\alpha\alpha$$
$$\varphi_{3} = |S = 1/2, M = 1/2\rangle = c_{1}\alpha\alpha\beta + c_{2}\alpha\beta\alpha + c_{3}\beta\alpha\alpha$$

Or in matrix form

$$\begin{pmatrix} \varphi_1 \\ \varphi_2 \\ \varphi_3 \end{pmatrix} = \begin{pmatrix} a_1 & a_2 & a_3 \\ b_1 & b_2 & b_3 \\ c_1 & c_2 & c_3 \end{pmatrix} \begin{pmatrix} \alpha \alpha \beta \\ \alpha \beta \alpha \\ \beta \alpha \alpha \end{pmatrix}$$

We may invert the matrix and write

$$\begin{pmatrix} \alpha \alpha \beta \\ \alpha \beta \alpha \\ \beta \alpha \alpha \end{pmatrix} = \begin{pmatrix} A_1 & A_2 & A_3 \\ B_1 & B_2 & B_3 \\ C_1 & C_2 & C_3 \end{pmatrix} \begin{pmatrix} \varphi_1 \\ \varphi_2 \\ \varphi_3 \end{pmatrix}$$

Where

$$\begin{pmatrix} A_1 & A_2 & A_3 \\ B_1 & B_2 & B_3 \\ C_1 & C_2 & C_3 \end{pmatrix} = \begin{pmatrix} a_1 & a_2 & a_3 \\ b_1 & b_2 & b_3 \\ c_1 & c_2 & c_3 \end{pmatrix}^{-1}$$

The spin function $\alpha\alpha\beta$ is a linear combination of the quartet, φ_1 and the two doublets, $\varphi_2 \& \varphi_3$

$$\alpha\alpha\beta = A_1\varphi_1 + A_2\varphi_2 + A_3\varphi_3$$

To determine φ_1 we will operate on $\alpha\alpha\beta$ with a projection operator that will annihilate $\varphi_2 \& \varphi_3$

$$\left(\hat{S}^{2} - \frac{1}{2}\left(\frac{1}{2} + 1\right)\right)\alpha\alpha\beta = A_{1}\left(\frac{3}{2}\left(\frac{3}{2} + 1\right) - \frac{1}{2}\left(\frac{1}{2} + 1\right)\right)\varphi_{1} = 3A_{1}\varphi_{1}$$

And since

$$\hat{S}^2 \alpha \alpha \beta = \left(\hat{S}_+ \hat{S}_- - S_z + S_z^2\right) \alpha \alpha \beta = \left(\hat{S}_+ \hat{S}_- - 1/4\right) \alpha \alpha \beta$$

And

$$\hat{S}_{+}\hat{S}_{-}\alpha\alpha\beta = \hat{S}_{+}\left(\beta\alpha\beta + \alpha\beta\beta\right) = 2\alpha\alpha\beta + \beta\alpha\alpha + \alpha\beta\alpha$$

So we have

$$\hat{S}^2 \alpha \alpha \beta = 7/4 \alpha \alpha \beta + \beta \alpha \alpha + \alpha \beta \alpha$$

And therefore

$$(\hat{S}^2 - 3/4)\alpha\alpha\beta = \alpha\alpha\beta + \beta\alpha\alpha + \alpha\beta\alpha = 3A_1\varphi_1$$

After normalization we have

$$\varphi_1 = |S = 3/2, M = 1/2\rangle = \frac{1}{\sqrt{3}}(\alpha\alpha\beta + \beta\alpha\alpha + \alpha\beta\alpha)$$

To find the doublets we annihilate the quartet.

$$\left(\hat{S}^{2} - \frac{3}{2}\left(\frac{3}{2} + 1\right)\right)\alpha\alpha\beta = \left(\frac{1}{2}\left(\frac{1}{2} + 1\right) - \frac{3}{2}\left(\frac{3}{2} + 1\right)\right)\left(A_{2}\varphi_{2} + A_{3}\varphi_{3}\right)$$

Note that since $\varphi_2 & \varphi_3$ are both doublets any linear combination is a doublet so

$$\left(\hat{S}^2 - \frac{3}{2}\left(\frac{3}{2} + 1\right)\right)\alpha\alpha\beta = -2\alpha\alpha\beta + \beta\alpha\alpha + \alpha\beta\alpha$$

is a doublet. The other linearly independent doublet is found by projecting the quartet out of one of the remaining functions, either $\alpha\beta\alpha$ or $\beta\alpha\alpha$.

$$\left(\hat{S}^2 - \frac{3}{2}\left(\frac{3}{2} + 1\right)\right)\alpha\beta\alpha = -2\alpha\beta\alpha + \beta\alpha\alpha + \alpha\alpha\beta \left(\hat{S}^2 - \frac{3}{2}\left(\frac{3}{2} + 1\right)\right)\alpha\alpha\beta = -2\alpha\alpha\beta + \beta\alpha\alpha + \alpha\beta\alpha$$

We then have three normalized doublet spin functions

$$\frac{1}{\sqrt{6}}(-2\alpha\alpha\beta + \beta\alpha\alpha + \alpha\beta\alpha)$$
$$\frac{1}{\sqrt{6}}(-2\alpha\beta\alpha + \beta\alpha\alpha + \alpha\alpha\beta)$$
$$\frac{1}{\sqrt{6}}(-2\beta\alpha\alpha + \alpha\beta\alpha + \alpha\alpha\beta)$$

Note that these three doublets can not be linearly independent because there are only two in the space. We can form two orthonormal functions by taking

$$\frac{1}{\sqrt{6}}(-2\alpha\alpha\beta+\beta\alpha\alpha+\alpha\beta\alpha)$$

as one doublet and forming the second by subtracting

$$\frac{1}{\sqrt{6}}(-2\alpha\beta\alpha + \beta\alpha\alpha + \alpha\alpha\beta)\,from\frac{1}{\sqrt{6}}(-2\beta\alpha\alpha + \alpha\beta\alpha + \alpha\alpha\beta)$$

We have, after normalization the second doublet.

$$\frac{1}{\sqrt{2}}(\alpha\beta\alpha-\beta\alpha\alpha).$$

We can form the remaining M = -1/2 functions by the same projection operator technique but it instructive to use the raising and lowering operators. Recall that for any angular momentum eigenfunction $|J,M\rangle$ operating with \hat{J}_{\pm} keeps J the same and changes M to $M \pm 1$

$$\hat{J}_{\pm} | J, M \rangle = A(J, M) | J, M \pm 1 \rangle$$

Given
$$\left|S=3/2, M=\frac{1}{2}\right\rangle$$
 we may form $\left|S=3/2, M=-\frac{1}{2}\right\rangle$ by operating with \hat{S}_{-} so

$$\left|S=3/2, M=-\frac{1}{2}\right\rangle \approx \hat{S}_{-}\left(\alpha\alpha\beta+\alpha\beta\alpha+\beta\alpha\alpha\right)=2\left(\alpha\beta\beta+\beta\alpha\beta+\beta\beta\alpha\right)$$

and after normalization we have

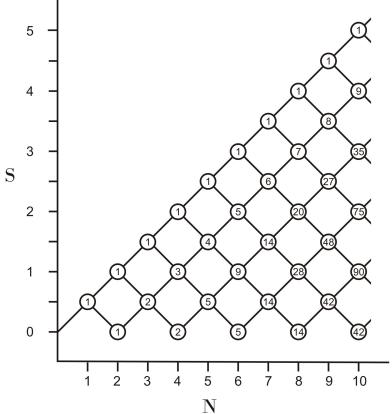
$$\left|S=3/2, M=-\frac{1}{2}\right\rangle = \frac{1}{\sqrt{3}}(\alpha\beta\beta+\beta\alpha\beta+\beta\beta\alpha)$$

Exercise

Show that the effect of the lowering operator \hat{S}_{-} on $\alpha\alpha\alpha$ is the same function $|S = 3/2, M = 1/2\rangle$ obtained by using the projection operator on $\alpha\alpha\beta$.

Branching Diagram

A very convenient way of keeping track of the number of spin eigenfunctions of a given multiplicity for an N electron system is the branching diagram shown below



The diagram is constructed as follows. A two-electron system is constructed from a one electron system by coupling the additional spin 1/2 and the original spin 1/2 into either a singlet (down) or a triplet (up). To form the spin eigenfunctions for a three electron system we couple the two electron triplet into a quartet (up) and a doublet (down). Additionally the singlet may be coupled (up) into a doublet. We then have 1 quartet and two doublets. To form the 4 electron system we couple the quartet into a quintet (up) and into a triplet (down). The two doublets couple up to form two additional triplets and down to form two singlets. The four electron system then can form 1 quintet (S = 2), 3 triplets (S = 1) and 2 singlets (S = 0) for a total of 16 states. This is of course the required 2⁴. Note the rapid increase in the

number of low spin eiegenfunctions as the number of unpaired electrons increases. For N=10 we have 42 singlets, 90 triplets, 75 quintets, 35 septets, 9 nonets and 1 bigtet.

We gather below a few explicit spin eigenfunction for N=1 to 5 and specifically for the case M = S. Lower values of M for a particular S may be generated using the lowering operator \hat{S}_- . These are from the Sanibel notes of Frank E Harris. Note that the functions have been factored where possible to highlight the physical couplings. For example the N=4 singlet, $(\alpha\beta - \beta\alpha)(\alpha\beta - \beta\alpha)/2$ obtains from the singlet coupling of electron pairs 1&2 and 3&4 while the N=5 doublet, $(\alpha\beta - \beta\alpha)(\alpha\beta - \beta\alpha)\alpha/2$ also obtains from the singlet coupling of electron pairs 1&2 and 3&4 with the doublet spin angular momentum being carried by the α spin of electron 5.

N	S	$\chi_{\scriptscriptstyle SS}$
1	1 / 2	
1	1/2	α
2	1	$(\alpha\beta + \beta\alpha)/\sqrt{2}$
2	0	$(\alpha\beta - \beta\alpha)/\sqrt{2}$
3	3/2	ααα
3	1/2	$(\alpha\beta - \beta\alpha)\alpha / \sqrt{2}$
3	1/2	$(\beta\alpha\alpha + \alpha\beta\alpha - 2\alpha\alpha\beta)/\sqrt{6}$
4	2	αααα
4	1	$(lphaeta - eta lpha) lpha lpha / \sqrt{2}$
4	1	$(\beta\alpha\alpha + \alpha\beta\alpha - 2\alpha\alpha\beta)\alpha / \sqrt{6}$
4	1	$(\beta\alpha\alpha\alpha + \alpha\beta\alpha\alpha + \alpha\alpha\beta\alpha - 3\alpha\alpha\alpha\beta)/\sqrt{12}$
4	0	$(\alpha\beta - \beta\alpha)(\alpha\beta - \beta\alpha)/2$
4	0	$(2\alpha\alpha\beta\beta+2\beta\beta\alpha\alpha-lphaetalphaeta-etalphaetalpha-lphaetaetaeta-etalphaetaeta/\sqrt{12}$
5	5/2	ααααα
5	3/2	$(\alpha\beta - \beta\alpha)\alpha\alpha\alpha / \sqrt{2}$
5	3/2	$(etalphalpha+lphaetalpha-2lphalphaeta)lphalpha/\sqrt{6}$
5	3/2	$(etalphalphalpha+lphaetalphalpha-3lphalphalpha)lpha/\sqrt{12}$
5	3/2	$(etalphalphalpha+lphaetalphalpha+lphalphaetaetalpha-4lphalphalphaeta)/\sqrt{20}$
5	1/2	$(\alpha\beta - \beta\alpha)(\alpha\beta - \beta\alpha)\alpha/2$
5	1/2	$(2\alpha\alpha\beta\beta+2\beta\beta\alpha\alpha-lpha\betalpha\beta-etalphaetalpha-lphaetaetaeta-etalphaetaeta-etalphalphaeta)lpha/\sqrt{12}$
5	1/2	$(\alpha\beta - \beta\alpha)(\beta\alpha\alpha + \alpha\beta\alpha - 2\alpha\alpha\beta)/\sqrt{12}$
5	1/2	$(\alpha\alpha\alpha\beta\beta + \alpha\alpha\beta\alpha\beta + \beta\beta\alpha\alpha\alpha - \alpha\beta\alpha\alpha\beta - \beta\alpha\alpha\alpha\beta - \alpha\alpha\beta\beta\alpha)/\sqrt{6}$
5	1/2	$(2\alpha\alpha\beta\alpha\beta - 2\alpha\alpha\alpha\beta\beta + \alpha\beta\alpha\beta\alpha + \beta\alpha\alpha\beta\alpha - \alpha\beta\beta\alpha\alpha - \beta\alpha\beta\alpha\alpha)/\sqrt{12}$

Orthonormal Spin Functions for N spins with total spin S through N=5