

1. (6 pts.) Let

$$H = \sum_{p,q} \langle p|z|q \rangle X_p^\dagger X_q + \frac{1}{2} \sum_{p,q,r,s} \langle pq|v|rs \rangle X_p^\dagger X_q^\dagger X_s X_r \quad (1)$$

be the Hamiltonian containing up to two-body terms. Use the anti-commutation relations for the creation and annihilation operators to show that if the domain of operator H is restricted to an N -electron Hilbert space with a fixed number of particles N , we can write

$$H = \frac{1}{2} \sum_{p,q,r,s} \langle pq|w|rs \rangle X_p^\dagger X_q^\dagger X_s X_r, \quad (2)$$

where

$$\langle pq|w|rs \rangle = \langle pq|v|rs \rangle + \frac{1}{N-1} (\delta_{pr} \langle q|z|s \rangle + \delta_{qs} \langle p|z|r \rangle). \quad (3)$$

2. (8 pts.) Let
- Z_N
- be a generic one-body operator in the normal-product form,

$$Z_N = \sum_{p,q} \langle p|z|q \rangle N[X_p^\dagger X_q], \quad (4)$$

and let T_1 be a one-body particle-hole excitation operator,

$$T_1 = \sum_{i,a} \langle a|t_1|i \rangle N[X_a^\dagger X_i]. \quad (5)$$

Use the hole-particle variant of Wick's theorem to prove that

$$[Z_N, T_1] |\Phi\rangle = \sum_{a,i} \left(\sum_b \langle a|z|b \rangle \langle b|t_1|i \rangle - \sum_j \langle a|t_1|j \rangle \langle j|z|i \rangle \right) |\Phi_i^a\rangle + \left(\sum_{a,i} \langle a|t_1|i \rangle \langle i|z|a \rangle \right) |\Phi\rangle, \quad (6)$$

where $[\ , \]$ is the commutator, $|\Phi\rangle$ is the Fermi vacuum state, and $|\Phi_i^a\rangle$ are the singly excited determinants relative to $|\Phi\rangle$. You will have a chance to prove the same identity using diagrams later. The above equation is a demonstration of the fact (used, for example, in coupled-cluster theory) that commutators of many-body operators produce *connected* quantities (diagrams) only.

3. (6 pts.) Use the results obtained in Problem 2, particularly Eq. (6), to derive the formulas for

$$\langle \Phi | [Z_N, T_1] | \Phi \rangle, \quad (7)$$

$$\langle \Phi_i^a | [Z_N, T_1] | \Phi \rangle, \quad (8)$$

and

$$\langle \Phi_{ij}^{ab} | [Z_N, T_1] | \Phi \rangle \quad (9)$$

in terms of matrix elements $\langle p|z|q \rangle$ and $\langle b|t_1|j \rangle$. As in Problem 2, $[\ , \]$ is the commutator, $|\Phi\rangle$ is the Fermi vacuum state, and $|\Phi_i^a\rangle$ is the singly excited determinant relative to

$|\Phi\rangle$. $|\Phi_{ij}^{ab}\rangle$ is the doubly excited determinant relative to $|\Phi\rangle$. The remaining definitions are the same as in Problem 2. In general, what are the values of n in the expression

$$\langle \Phi_{i_1 \dots i_n}^{a_1 \dots a_n} | [Z_N, T_1] | \Phi \rangle, \quad (10)$$

where $|\Phi_{i_1 \dots i_n}^{a_1 \dots a_n}\rangle$ is the n -tuply excited determinant relative to $|\Phi\rangle$, that produce non-zero results? Justify your answer. How would your answer change if we considered

$$\langle \Phi_{i_1 \dots i_n}^{a_1 \dots a_n} | Z_N T_1 | \Phi \rangle \quad (11)$$

instead (i.e., if we replaced the commutator of Z_N and T_1 by their product)?