

CONTINUOUS SPECTRA (ONLY)

$$\checkmark \quad \hat{F} |f_\mu\rangle = \lambda_\mu |f_\mu\rangle$$

Annotations:
- "set λ_μ " points to λ_μ
- "solve for $|f_\mu\rangle$ for a given λ_μ " points to $|f_\mu\rangle$

μ is a continuous variable

$\lambda_\mu \equiv \lambda(\mu)$ is a function of μ .

$f_\mu(q) \equiv f(q; \mu)$ is

$q_1, \dots, q_f$
a function of q parameterized by μ .

$\mu \rightarrow \mu + \Delta\mu, \lambda_\mu \rightarrow \lambda_{\mu+\Delta\mu}$
 $|f_{\mu+\Delta\mu}\rangle$ or $f_{\mu+\Delta\mu}(q)$
 is a solution corresponding
 to $\lambda_{\mu+\Delta\mu}$.

Example:

$$\hat{F} = \hat{p}_x = \frac{\hbar}{i} \frac{\partial}{\partial x}$$

$$\hat{p}_x u_k(x) = \hbar k u_k(x)$$

$$u_k(x) = A e^{ikx}, \quad \int_{-\infty}^{\infty} |u_k(x)|^2 dx = \infty$$

Conventional normalization
does not work.

We will prove that $|f_\mu\rangle$ is
orthogonal to $|f_\nu\rangle$ if $\mu \neq \nu$.

We start from the statement
that all $|f_\mu\rangle$ s form a
complete set,

$$|\varphi\rangle = \int c_\mu |f_\mu\rangle d\mu \quad \checkmark$$

$$\varphi(q) = \int c_\mu f_\mu(q) d\mu$$

In a discrete case (assuming $\|\varphi\|=1$)

$$|c_\mu|^2 = P(F = \lambda_\mu)$$

In a continuous case

$|c_\mu|^2 d\mu$ = probability
that F has a value from
the interval $(\lambda_\mu, \lambda_\mu + d\lambda_\mu)$

$$\checkmark |c_\mu|^2 d\mu = P[FE(\lambda_\mu, \lambda_{\mu+d\lambda_\mu})]$$

Assuming that $\|\varphi\| = 1$,

$$\int \varphi^*(q) \varphi(q) d\tau = 1$$

$$\int |c_\mu|^2 d\mu = 1 \quad \checkmark$$

$$P[FE(\lambda_\mu, \lambda_{\mu+d\lambda_\mu})]$$

$$1 = \langle \varphi | \varphi \rangle = \left\langle \int c_\mu f_\mu d\mu \middle| \varphi \right\rangle$$

$$= \int \underline{c_\mu}^* \langle f_\mu | \varphi \rangle d\mu \quad \left. \vphantom{\int} \right\}$$

$$= \int \underline{c_\mu}^* c_\mu d\mu$$

$$\int c_{\mu}^* [\overbrace{\langle f_{\mu} | \varphi \rangle - c_{\mu}}^0] d\mu = 0$$

$$\underline{c_{\mu} = \langle f_{\mu} | \varphi \rangle} \quad (\text{as in discrete case})$$

Orthogonality of $|f_{\mu}\rangle$ s:

$$\begin{aligned} c_{\mu} &= \langle f_{\mu} | \varphi \rangle \\ &= \langle f_{\mu} | \int c_{\nu} f_{\nu} d\nu \rangle \\ &= \int c_{\nu} \langle f_{\mu} | f_{\nu} \rangle d\nu \end{aligned}$$

Compare this with

$$\checkmark \quad \int f(x) \delta(x-y) dx = f(y)$$

$$\forall c(\mu) = \int c(\nu) g(\mu, \nu) d\nu$$

$$g(\mu, \nu) = \langle f_\mu | f_\nu \rangle$$

$$g(\mu, \nu) = \delta(\mu - \nu)$$

$$\int c(\nu) \delta(\mu - \nu) d\nu = c(\mu)$$

$$\langle f_\mu | f_\nu \rangle = \delta(\mu - \nu)$$

$$\text{If } \mu \neq \nu, \langle f_\mu | f_\nu \rangle = 0$$

$$|f_\mu\rangle \perp |f_\nu\rangle$$

(again similar to discrete case)

Closure: $\langle f_\mu | \varphi \rangle$

$$\varphi(q) = \int c_\mu f_\mu(q) d\mu$$

$$\langle f_\mu | \varphi \rangle = \int f_\mu^*(q') \varphi(q') d\tau'$$

$$\varphi(q) = \int d\mu \int d\tau' f_\mu^*(q') \varphi(q') f_\mu(q)$$

$$= \int d\tau' \varphi(q') \int d\mu f_\mu^*(q') f_\mu(q)$$

$$\varphi(q) = \int d\tau' \varphi(q') \left[\int f_\mu^*(q') f_\mu(q) d\mu \right]$$

$$\varphi(q) = \int d\tau' \varphi(q') \delta(q' - q) d\tau'$$

$$\delta(q' - q) = \prod_{l=1}^f \delta(q'_l - q_l)$$

$$\int f_{\mu}^{*}(q') f_{\mu}(q) d\mu = \delta(q' - q)$$

(in a discrete case,

$$\sum_{\mu} f_{\mu}^{*}(q') f_{\mu}(q) = \delta(q' - q)$$

Parseval's relation:

$$\langle \varphi | \gamma \rangle = \int \varphi^{*}(q) \gamma(q) dq$$

$$= \int \int \varphi^{*}(q) \underbrace{\delta(q' - q)} \gamma(q') dq' dq$$

$$= \int \int \underbrace{\varphi^*(q)}_{\text{pink}} \left[\int \underbrace{f_\mu^*(q')}_{\text{green}} \underbrace{f_\mu(q)}_{\text{pink}} d\mu \right] \\ \times \underbrace{\psi(q')}_{\text{green}} \underbrace{d\tau'}_{\text{green}} \underbrace{d\tau}_{\text{pink}}$$

$$= \int d\mu \left[\int \underbrace{\varphi^*(q)}_{\text{pink}} \underbrace{f_\mu(q)}_{\text{pink}} d\tau \right] \\ \times \left[\int \underbrace{f_\mu^*(q')}_{\text{green}} \underbrace{\psi(q')}_{\text{green}} d\tau' \right]$$

$$= \int d\mu \langle \varphi | f_\mu \rangle \langle f_\mu | \psi \rangle$$

$$\langle \varphi | \psi \rangle = \int \langle \varphi | \underbrace{f_\mu}_{\text{orange}} \rangle \langle \underbrace{f_\mu}_{\text{orange}} | \psi \rangle d\mu$$

(in a discrete case,

$$\langle \phi | \psi \rangle = \sum_{\mu} \langle \phi | f_{\mu} \rangle \langle f_{\mu} | \psi \rangle$$

$$\int |f_{\mu} \rangle \langle f_{\mu}| d\mu = 1$$

Expectation value:

In a discrete case

$$\begin{aligned} \sum_{\mu} \lambda_{\mu} P(F = \lambda_{\mu}) & \quad \checkmark \\ = \sum_{\mu} \lambda_{\mu} |c_{\mu}|^2 & = \dots = \\ = \langle \phi | \hat{F} | \phi \rangle \end{aligned}$$

Continuous case:

$$\langle \hat{F} \rangle = \int \lambda_{\mu} P[FE(\lambda_{\mu}, \lambda_{\mu} + d\lambda_{\mu})]$$

$|c_{\mu}|^2 d\mu$

$$= \int \lambda_{\mu} |c_{\mu}|^2 d\mu$$

$$= \int \lambda_{\mu} c_{\mu}^* c_{\mu} d\mu$$

$$c_{\mu} = \langle f_{\mu} | \varphi \rangle$$

$$= \int \lambda_{\mu} \langle f_{\mu} | \varphi \rangle^* \langle f_{\mu} | \varphi \rangle d\mu$$

$$\stackrel{(a)}{=} \int \lambda_{\mu} \langle \varphi | f_{\mu} \rangle \langle f_{\mu} | \varphi \rangle d\mu$$

$$= \int \langle \varphi | \underbrace{\hat{F} f_\mu}_{\hat{F} f_\mu} \rangle \langle f_\mu | \varphi \rangle d\mu$$

$$= \int \langle \varphi | \hat{F} f_\mu \rangle \langle f_\mu | \varphi \rangle d\mu$$

TOR

$$= \int \langle \hat{F} \varphi | f_\mu \rangle \langle f_\mu | \varphi \rangle d\mu$$

closure

$$\langle \hat{F} \varphi | \varphi \rangle \stackrel{\text{TOR}}{=} \langle \varphi | \hat{F} \varphi \rangle$$

$$\langle \hat{F} \rangle = \langle \varphi | \hat{F} | \varphi \rangle$$

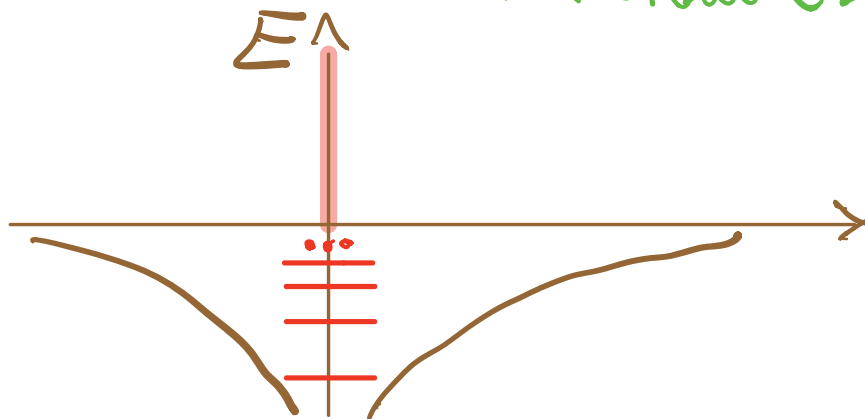
$$(\|\varphi\|=1)$$

GENERAL CASE (DISCRETE & CONTINUOUS SPECTRA)

$$\hat{F}|\psi_\mu\rangle = \lambda_\mu |\psi_\mu\rangle$$

In a discrete case, μ has discrete values, λ_μ s are eigenvalues.

In a continuous case, μ is a continuous variable.



$\{ |f_\mu\rangle \}$ is complete:

$$| \varphi \rangle = \sum_{\mu} c_{\mu} |f_{\mu}\rangle \quad \leftarrow \text{sum over } \mu$$

(discrete part)

$$+ \int c_{\mu} |f_{\mu}\rangle d\mu \quad \leftarrow \text{integrate over } \mu$$

(continuous part)

$$c_{\mu} = \langle f_{\mu} | \varphi \rangle.$$

New symbol

$$\sum_{\mu} \begin{cases} \nearrow \sum_{\mu} \text{ in a discrete case} \\ \searrow \int d\mu \text{ in a continuous case} \end{cases}$$

$$|\varphi\rangle = \sum_{\mu} c_{\mu} |f_{\mu}\rangle$$

$$\sum_{\mu} |c_{\mu}|^2 + \int |c_{\mu}|^2 d\mu = 1 \quad \text{if } \|\varphi\|=1$$

$$\sum_{\mu} |c_{\mu}|^2 = 1$$

Closure:

$$\sum_{\mu} f_{\mu}^*(q') f_{\mu}(q) + \int f_{\mu}^*(q') f_{\mu}(q) d\mu$$

$$= \sum_{\mu} f_{\mu}^{*}(q') f_{\mu}(q)$$

$$= \delta(q' - q) -$$

In a discrete part of the spectrum,

$$\langle f_{\mu} | f_{\nu} \rangle = \delta_{\mu\nu}.$$

In a continuous part,

$$\langle f_{\mu} | f_{\nu} \rangle = \delta(\mu - \nu).$$

Resolution of identity,

$$\sum_{\mu} |f_{\mu}\rangle \langle f_{\mu}| + \int |f_{\mu}\rangle \langle f_{\mu}| d\mu$$

$$= 1 = \sum_{\mu} |f_{\mu}\rangle \langle f_{\mu}|$$

Spectral theorem:

$$\hat{F} = \sum_{\mu} \lambda_{\mu} |f_{\mu}\rangle \langle f_{\mu}| + \sum_{\mu} \lambda_{\mu}^* |f_{\mu}\rangle \langle f_{\mu}|$$

$$= \sum_{\mu} \lambda_{\mu} |f_{\mu}\rangle \langle f_{\mu}|$$

$$\hat{F} = \sum_{\mu} \lambda_{\mu} P_{\langle f_{\mu} \rangle}$$

Expectation value: discrete

$$\langle \hat{F} \rangle = \sum_{\mu} \lambda_{\mu} |c_{\mu}|^2 + \int \lambda_{\mu} |c_{\mu}|^2 d\mu \quad (||\varphi||=1)$$

continuous



$$\langle \hat{F} \rangle = \langle \varphi | \hat{F} | \varphi \rangle.$$

Example of a continuous spectrum: momentum operator in 3D

$$\hat{\vec{p}} = (\hat{p}_x, \hat{p}_y, \hat{p}_z)$$

$$\hat{p}_x = \frac{\hbar}{i} \frac{\partial}{\partial x}$$

$$\hat{p}_y = \frac{\hbar}{i} \frac{\partial}{\partial y}$$

$$\hat{p}_z = \frac{\hbar}{i} \frac{\partial}{\partial z}$$

$$(\hat{\vec{p}} = \frac{\hbar}{i} \vec{\nabla})$$

$$\hat{\vec{p}} u_{\vec{p}} = \vec{p} u_{\vec{p}}$$

$$\frac{\hbar}{i} \frac{\partial}{\partial x} u_{\vec{p}}(x, y, z)$$

$$= p_x u_{\vec{p}}(x, y, z)$$

$$(x, y, z) = \vec{r}$$

$$\frac{\hbar}{i} \frac{\partial}{\partial y} u_{\vec{p}}(\vec{r}) = \underbrace{p_y}_{\text{wavy}} u_{\vec{p}}(\vec{r})$$

$$\frac{\hbar}{i} \frac{\partial}{\partial z} u_{\vec{p}}(\vec{r}) = \underbrace{p_z}_{\text{wavy}} u_{\vec{p}}(\vec{r})$$

$$u_{\vec{p}}(\vec{r}) = C e^{i(\vec{p} \cdot \vec{r})/\hbar}$$

$$= C e^{i(p_x x + p_y y + p_z z)/\hbar}$$

$$\frac{\vec{p}}{\hbar} = \vec{k} \quad \checkmark$$

$$u_{\vec{p}}(\vec{r}) = C e^{\underbrace{i \vec{k} \cdot \vec{r}}_{\text{wavy}}}$$

$$= C e^{i(\underbrace{k_x}_{\text{wavy}} x + \underbrace{k_y}_{\text{wavy}} y + \underbrace{k_z}_{\text{wavy}} z)}$$

✓ (in 1D, $\hat{p}_x u_p(x) = p_x u_p(x)$)

$$u_{\vec{k}}(\vec{r}) = C e^{i\vec{k} \cdot \vec{r}} \quad \checkmark$$

$$\hat{\vec{p}} u_{\vec{k}}(\vec{r}) = (\hbar \vec{k}) u_{\vec{k}}(\vec{r})$$

$$\langle f_\mu | f_\nu \rangle = \delta(\mu - \nu)$$

$$\langle u_{\vec{l}} | u_{\vec{k}} \rangle = \int_{\mathbb{R}^3} \underbrace{u_{\vec{l}}(\vec{r})^*}_{\text{green wavy}} \underbrace{u_{\vec{k}}(\vec{r})}_{\text{green wavy}} dx dy dz$$

$$= |C|^2 \int_{-\infty}^{\infty} e^{i(\underbrace{k_x}_{\text{green wavy}} - \underbrace{l_x}_{\text{purple wavy}})x} dx$$

$\nwarrow 2\pi \delta(k_x - l_x)$

$$\times \int_{-\infty}^{\infty} e^{i(\underbrace{k_y}_{\text{green wavy}} - \underbrace{l_y}_{\text{purple wavy}})y} dy$$

$\nwarrow 2\pi \delta(k_y - l_y)$

$$\times \int_{-\infty}^{\infty} e^{i(\underbrace{k_z}_{\text{green}} - \underbrace{l_z}_{\text{purple}})z} dz \quad \leftarrow 2\pi \delta(k_z - l_z)$$

$$u_{\vec{k}}(\vec{r}) = C e^{ik_x x} e^{ik_y y} \times e^{ik_z z}$$

$$\checkmark \int_{-\infty}^{\infty} e^{iax} dx = 2\pi \delta(a)$$

$$\langle u_{\vec{l}} | u_{\vec{k}} \rangle = (2\pi)^3 |C|^2$$

$$\times \delta(k_x - l_x) \delta(k_y - l_y) \times \delta(k_z - l_z)$$

$$\parallel \delta(\vec{k} - \vec{l})$$

$$\langle u_{\vec{l}} | u_{\vec{k}} \rangle = (2\pi)^3 |C|^2$$

$$\times \delta(\vec{k} - \vec{l})$$

$$= \delta(\vec{k} - \vec{l})$$

$$(2\pi)^3 |C|^2 = 1$$

$$C = \frac{1}{(2\pi)^{3/2}}$$

$$u_{\vec{k}}(\vec{r}) = \frac{1}{(2\pi)^{3/2}} e^{i\vec{k} \cdot \vec{r}}$$

$$\langle u_{\vec{l}} | u_{\vec{k}} \rangle = \delta(\vec{k} - \vec{l})$$

In 1D,



$$\hat{p}_x u_k(x) = \hbar k u_k(x)$$

$$u_k(x) = \frac{1}{\sqrt{2\pi}} e^{ikx}$$

$$\langle u_l | u_k \rangle = \delta(k-l).$$

Closure:

$$\int f_{\mu}^*(q') f_{\mu}(q) d\mu = \delta(q' - q)$$

$$\int u_{\vec{k}}^*(\vec{r}') u_{\vec{k}}(\vec{r}) dk_x dk_y dk_z$$


$$\begin{aligned}
&= \frac{1}{(2\pi)^3} \int_{-\infty}^{\infty} e^{ik_x(x-x')} dk_x \\
&\quad \times \int_{-\infty}^{\infty} e^{ik_y(y-y')} dk_y \\
&\quad \times \int_{-\infty}^{\infty} e^{ik_z(z-z')} dk_z \\
&\quad \int_{-\infty}^{\infty} e^{iax} dx = 2\pi \delta(a)
\end{aligned}$$

$$\begin{aligned}
&= \frac{1}{(2\pi)^3} 2\pi \delta(x-x') \\
&\quad \times 2\pi \delta(y-y')
\end{aligned}$$

$$\times \cancel{2\pi} \delta(z-z')$$

$$= \delta(x-x') \delta(y-y') \\ \times \delta(z-z')$$

$$= \delta(\vec{r}-\vec{r}')$$

$$\int u_{\vec{k}}(\vec{r}')^* u_{\vec{k}}(\vec{r}) dk_x dk_y dk_z \\ = \delta(\vec{r}-\vec{r}')$$

In 1D,

$$\int_{-\infty}^{\infty} u_k(x')^* u_k(x) dk \\ = \delta(x'-x)$$