

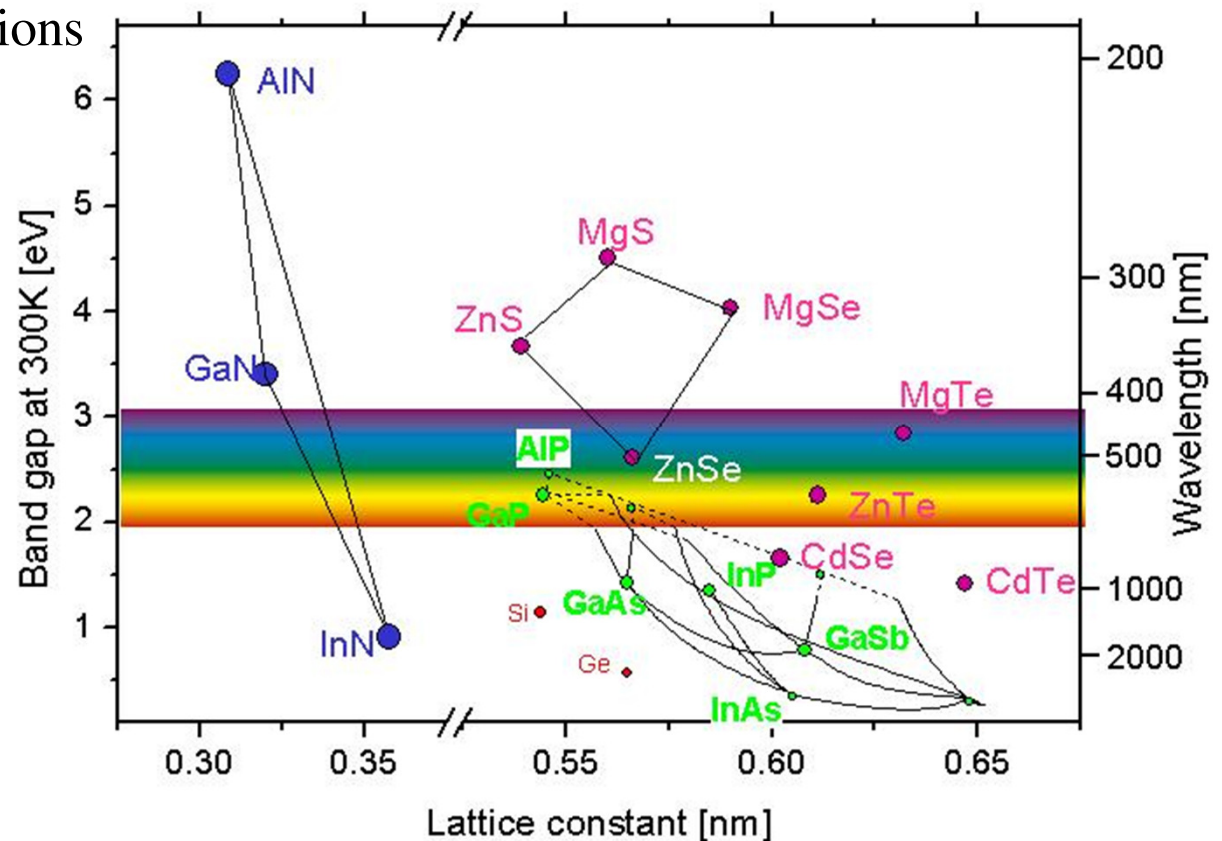
# Week 8: Ch. 11 Semiconductor diodes

## Principles of Scintillation Counters

### Semiconductor Diodes

- basics of semiconductors
- pure elements & dopants
- 5-3 Materials
- ion collection, leakage current
- diode structure, pn, np junctions
- depletion region

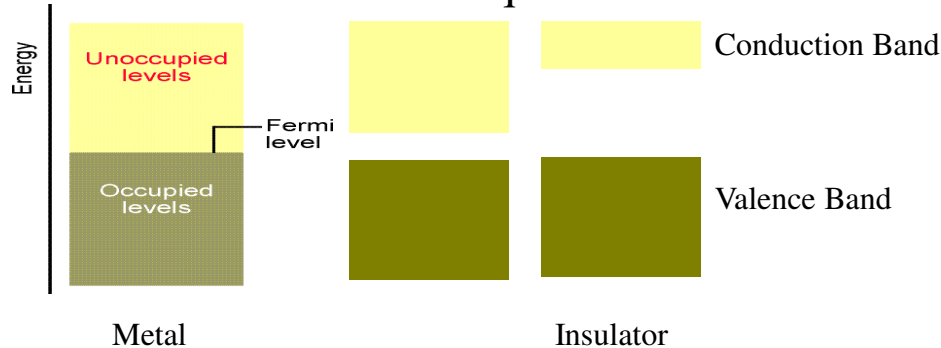
### Germanium Diodes



# Chap. 11 – Semiconductor Diodes

Semiconductor diodes provide the best resolution for energy measurements, silicon based devices are generally used for charged-particles, germanium for photons.

- Scintillators require  $\sim 100$  eV / “information carrier” .. Photoelectrons in this case
- Gas counters require  $\sim 35$  eV / “information carrier” .. Ion-pairs or Electrons
- Solid-state devices require  $\sim 3$  eV / “information carrier” .. Electron/hole pairs



A semiconductor is an insulator with a small band gap,  $\sim 1.2$  eV for silicon. Generally want smallest band gap *but* thermal excitation across the gap provides a leakage current. N.B. the actual band gap depends on the direction relative to the lattice (Si and Ge do not crystallize in cubic lattices) and the gap decreases slowly with temperature.

The ratio of ‘w’ to band gap is approximately constant for a wide range of materials – division of excitation energy between e/h pair and phonons, etc. is  $\sim$  constant.

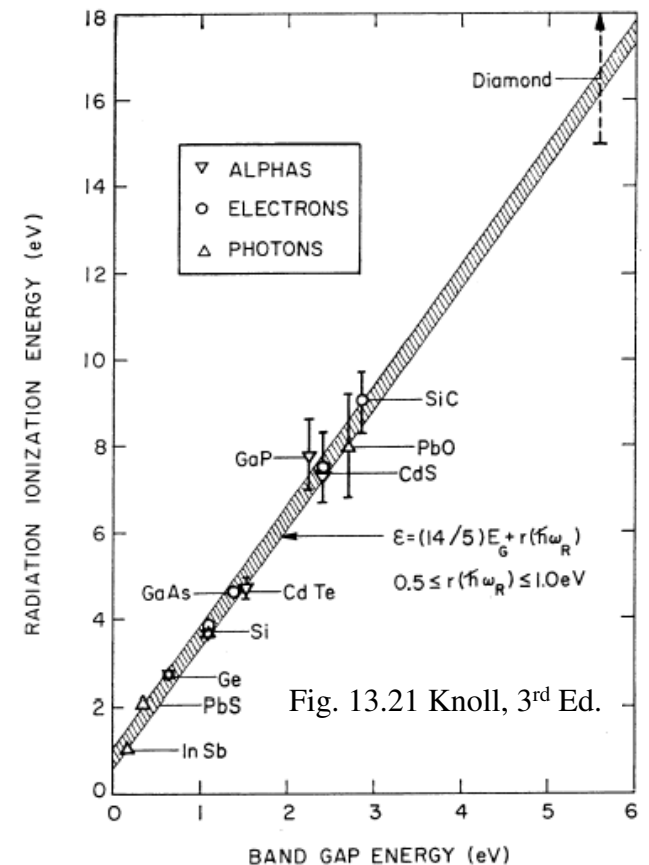


Fig. 13.21 Knoll, 3<sup>rd</sup> Ed.

# Semiconductors – Charge carriers

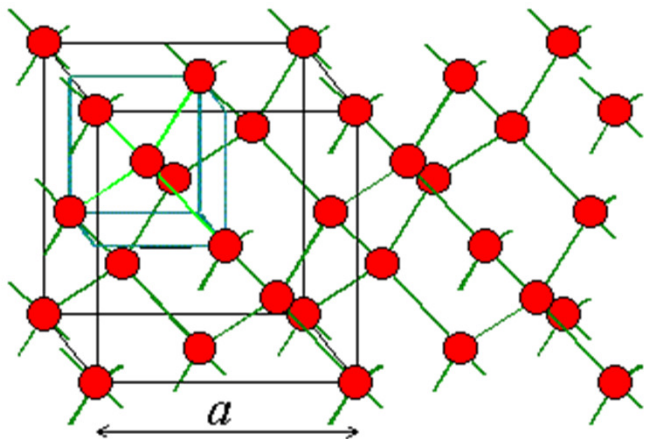
“pure” material, no dopants is called “intrinsic”

The intrinsic carrier density in a semiconductor is low:

$$\rho_e \sim \sqrt{N_V N_C} e^{-\varepsilon/2k_B T} \quad k_B T = 0.026 \text{ eV @ } 25^\circ \text{ C}$$

$$\rho_e \sim \sqrt{10^{19} 10^{19}} e^{-20} \sim 10^9 - 10^{10} \text{ cm}^{-3}$$

$N_V$  and  $N_C$  are the densities of states in the valence and conduction bands. ( Only rough estimates given here. )



“diamond lattice”

Lattice Constant

- Carbon      0.356 nm
- Silicon     0.543 nm
- Germanium 0.565 nm

14	
12.0111	-4
<b>C</b>	+2
6	+4
2-4	
28.0855	-4
<b>Si</b>	+2
14	+4
2-8-4	
72.59	-4
<b>Ge</b>	+2
32	+4
2-8-18-4	

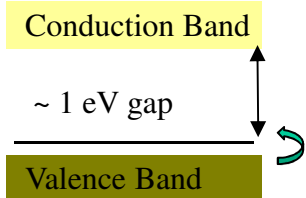
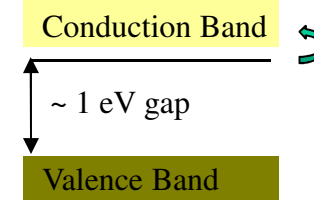
# Semiconductors – Dopants

13	14	15
10.81 <b>B</b> 5 2-3	12.0111 <b>C</b> 6 2-4	14.0067 <b>N</b> 7 2-5
26.98154 <b>Al</b> 13 2-8-3	28.0855 <b>Si</b> 14 2-8-4	30.97376 <b>P</b> 15 2-8-5
69.72 <b>Ga</b> 31 2-8-18-3	72.59 <b>Ge</b> 32 2-8-18-4	74.9216 <b>As</b> 33 2-8-18-5

- Add atoms from the neighboring groups in the periodic table
- Group 15, Phosphorous, nearly same size, excess electron
  - Group 13, Boron, nearly same size, electron deficit

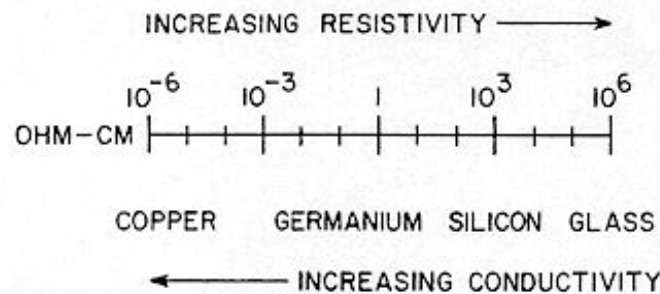
Donor level from P atom below conduction band by ~ 0.05 eV ..  
Thermally excite from donor, excess electrons → n-type

$$e^{-0.05/2kT} \sim e^{-1}$$



Acceptor level from B atom above valence band by ~ 0.05 eV ..  
Thermally excited from valence band, excess holes → p-type

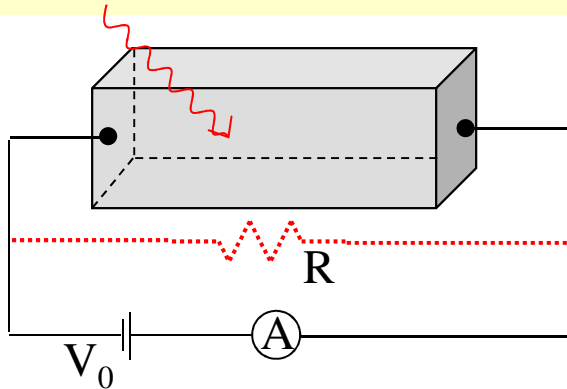
Control the conductivity by controlling the amount of dopants!  
N.B. 2 ppb gives  $(2 \times 10^{-9}) (5 \times 10^{22} / \text{cm}^3) = 10^{14} / \text{cm}^3 \gg 10^9$  for Si



For an n-type material, the electrons carry the current so that the resistivity is:

$$\rho = \frac{1}{q_e N_D \mu_e}$$

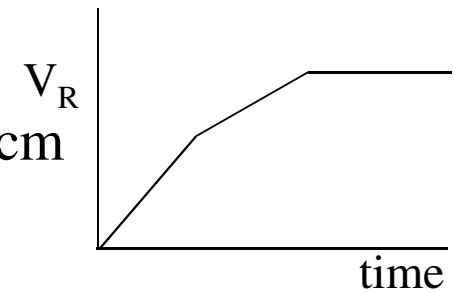
# Semiconductor – Ion Chamber?



Imagine constructing a simple block of intrinsic semiconductor and trying to use it as an ion chamber ... The block has a length, “L” and a cross sectional area, “A” with a resistivity of  $\rho = 60\text{k ohm-cm}$  (high quality silicon). Apply nominal  $V_0 = 60\text{V}$  bias to collect ions.

1 MeV energy into material creates  $\sim 3 \times 10^5$  e/h in  $\sim 50$  ns ... limiting drift velocity  $\sim 10^7$  cm/s

$$i_{\text{signal}} \sim \frac{\Delta q}{\Delta t} = \frac{3 \times 10^5 (1.6 \times 10^{-19})}{(L \text{ cm} / 10^7 \text{ cm} / \text{s})} = \frac{5}{L} 10^{-7} \text{ Amps for } L \text{ in cm}$$

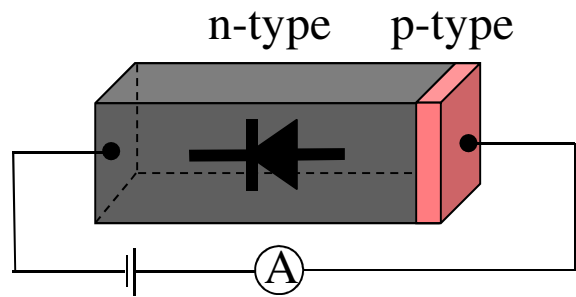


$$I_{\text{Leakage}} = \frac{V_0}{R} \quad \text{where} \quad R = \rho \frac{L}{A}$$

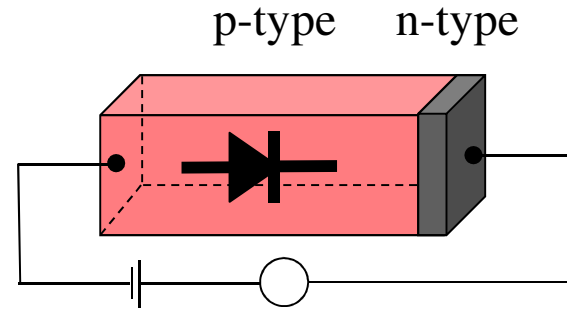
$$I_{\text{Leakage}} = \frac{V_0 A}{\rho L} \quad \rightarrow \quad \frac{60A}{60,000L} = \frac{A}{L} 10^{-3} \text{ Amps for } A/L \text{ in cm}$$

Thus,  $I \gg i$  so we need a trick to kill the leakage current.

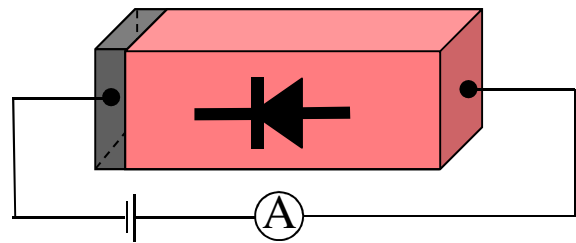
# Semiconductor Diodes – 1



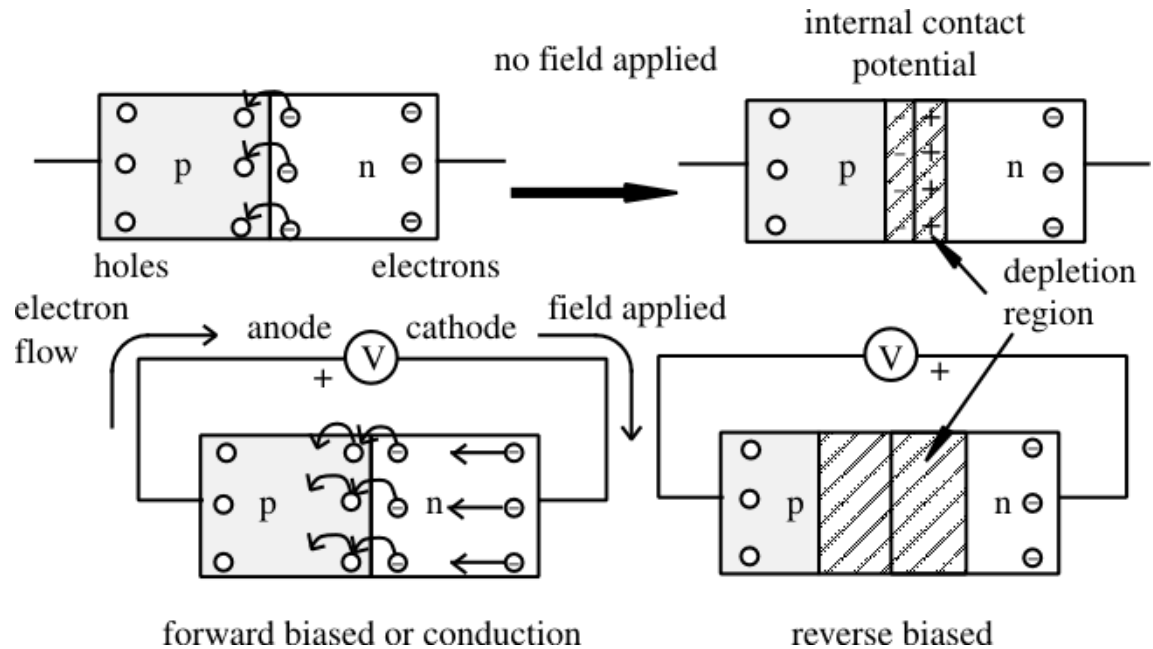
“Forward” bias – normal flow of current



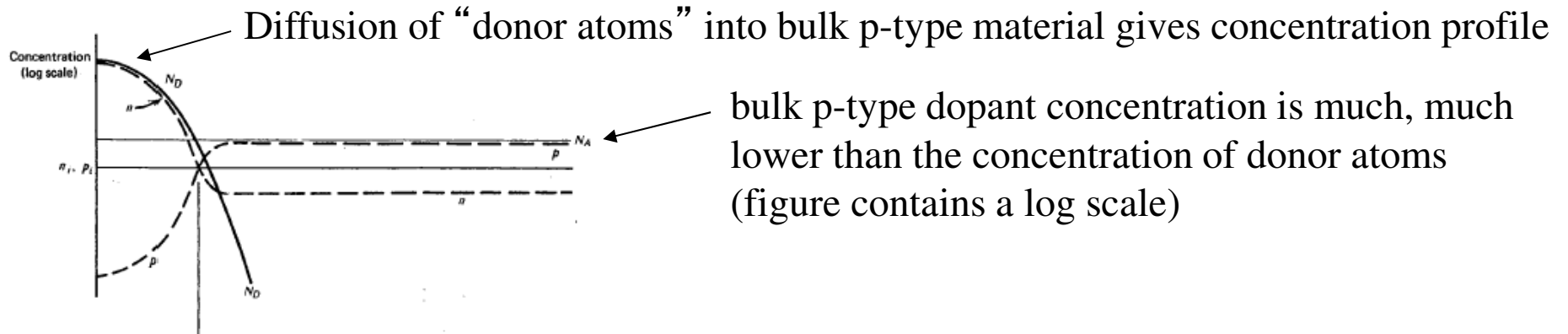
“Reverse” bias – no nominal current



The “type” is determined by the implanted atoms .. Different atoms can be put into a single piece of semiconductor, then an internal field will form due to migration of charges.



# Semiconductor Diodes – 2



Charge density,  $\rho(x)$ : Migration of the charge across the boundary causes a charge separation and a “depletion region” ... in this case the holes migrate for a longer distance due to imbalance in concentrations in two regions

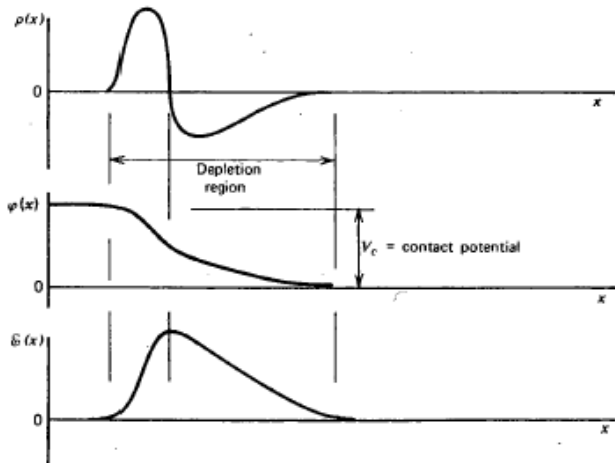


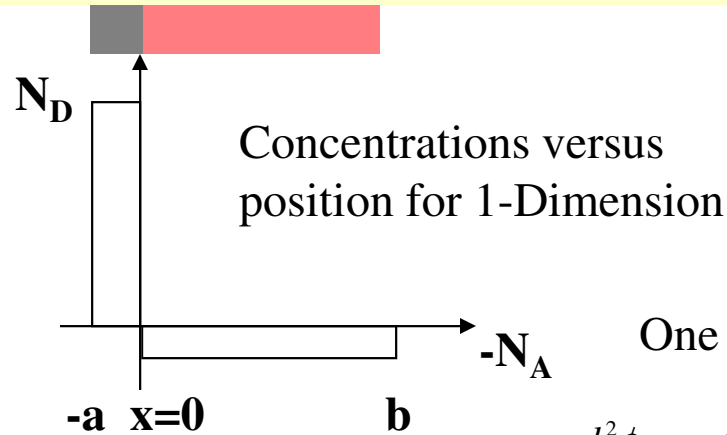
Fig. 11.8 Knoll, 3<sup>rd</sup>, 4<sup>th</sup> Eds.

Which creates an internal electric potential,  $\phi(x)$  with a potential difference of  $\sim 1V$

and an electric field,  $E(x)$ , where the lines of force originate on positive ions and terminate on the negative ions.

The depleted region has a very low concentration of mobile charge carriers and a very high resistivity – this is a very good region to measure/collect ionization.

# Semiconductor Diodes – Model for Depletion Depth



$$\text{Charge neutrality gives: } a \cdot N_D = b \cdot N_A \quad (1)$$

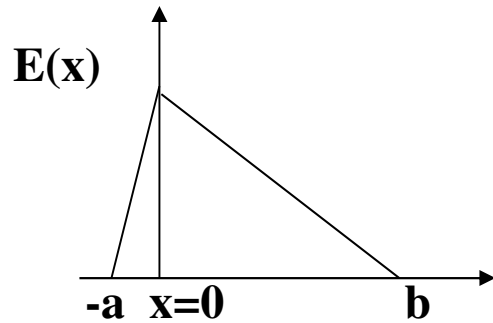
Depletion depth,  $d = a + b$

One Dimensional Poisson Equation in two regions:

$$\frac{d^2\phi}{dx^2} = \frac{-\rho(x)}{\epsilon} = \begin{cases} \frac{-q_e N_D}{\epsilon} & a < x \leq 0 \\ \frac{+q_e N_A}{\epsilon} & 0 < x \leq b \end{cases}$$

Boundary Conditions:

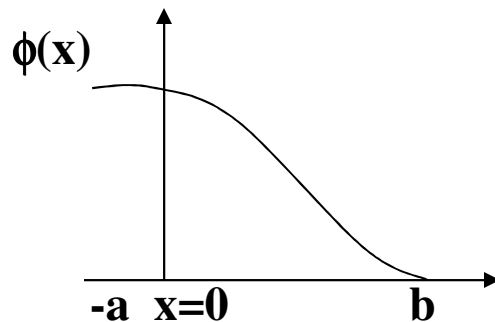
$$E(x = -a) = 0 \text{ \& } E(x = b) = 0$$



$$-E(x) = \frac{d\phi}{dx} = \begin{cases} \frac{-q_e N_D}{\epsilon} (x + a) & a < x \leq 0 \\ \frac{+q_e N_A}{\epsilon} (x - b) & 0 < x \leq b \end{cases}$$

Boundary Conditions:

$$\phi(x = -a) = V \text{ \& } \phi(x = b) = 0$$



$$\phi(x) = \begin{cases} \frac{-q_e N_D}{2\epsilon} (x + a)^2 + V & a < x \leq 0 \\ \frac{+q_e N_A}{2\epsilon} (x - b)^2 & 0 < x \leq b \end{cases}$$

Match at  $x = 0$ :

$$\frac{-q_e N_D}{2\epsilon} a^2 + V = \frac{+q_e N_A}{2\epsilon} b^2 \quad (2)$$

Depletion depth: combine (1) & (2)  $(a+b) \cdot b = 2\epsilon V / q_e N_A$

Generally  $b > a \dots d = b \sim (2\epsilon V / q_e N_A)^{1/2}$



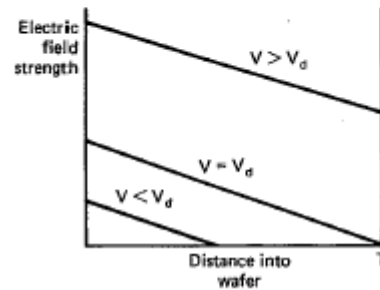
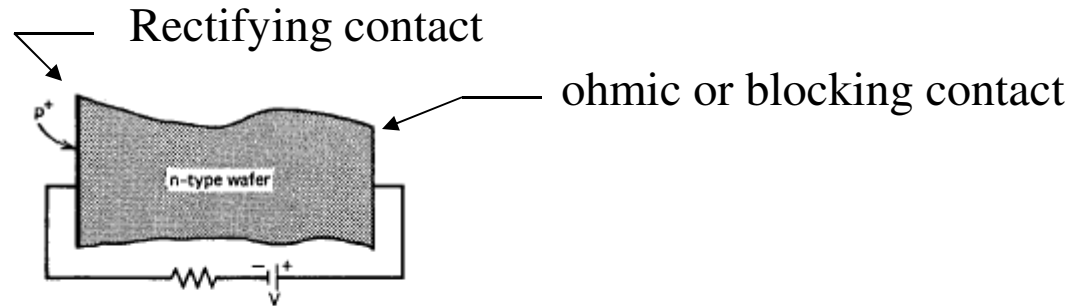
# Semiconductor Diodes – 3

Some properties of the junction:

$$d \sim \sqrt{\frac{2\epsilon V}{q_e N_A}} \quad \& \quad \rho = \frac{1}{q_e N_D \mu_e}$$

$$d \sim \sqrt{2\epsilon V \rho \mu} \quad \& \quad C = \frac{\epsilon A}{d}$$

$$E = \Delta V / \Delta x \quad \rightarrow \quad E_{\max} = \frac{2V}{d}$$



One can apply an external (reverse) bias voltage such that the depletion layer extends from the junction to the full thickness of the device

Fig. 11.12 Knoll, 3<sup>rd</sup>, 4<sup>th</sup> Eds.

## Silicon Surface Barrier Device

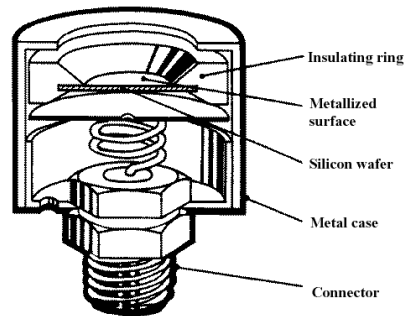
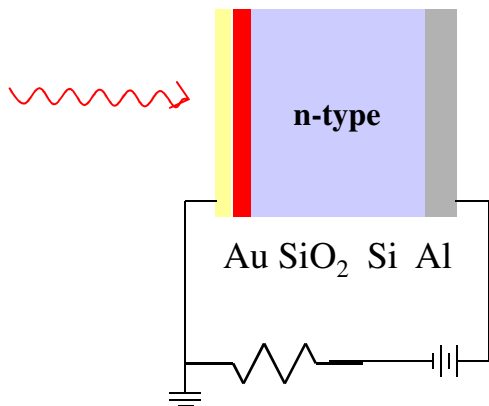


Fig. 11.11 Knoll, 3<sup>rd</sup>, 4<sup>th</sup> Eds. (From EG&G Ortec)

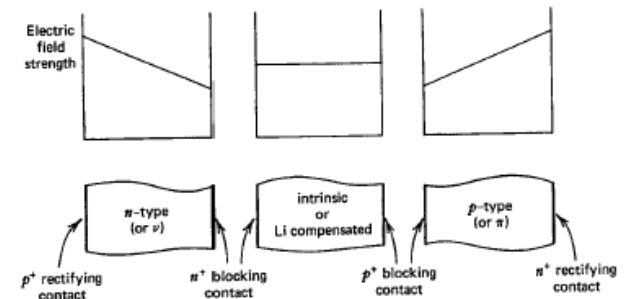
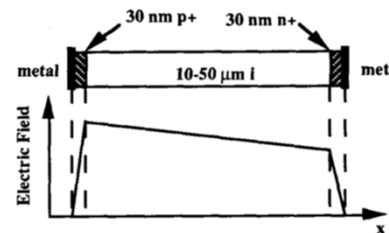
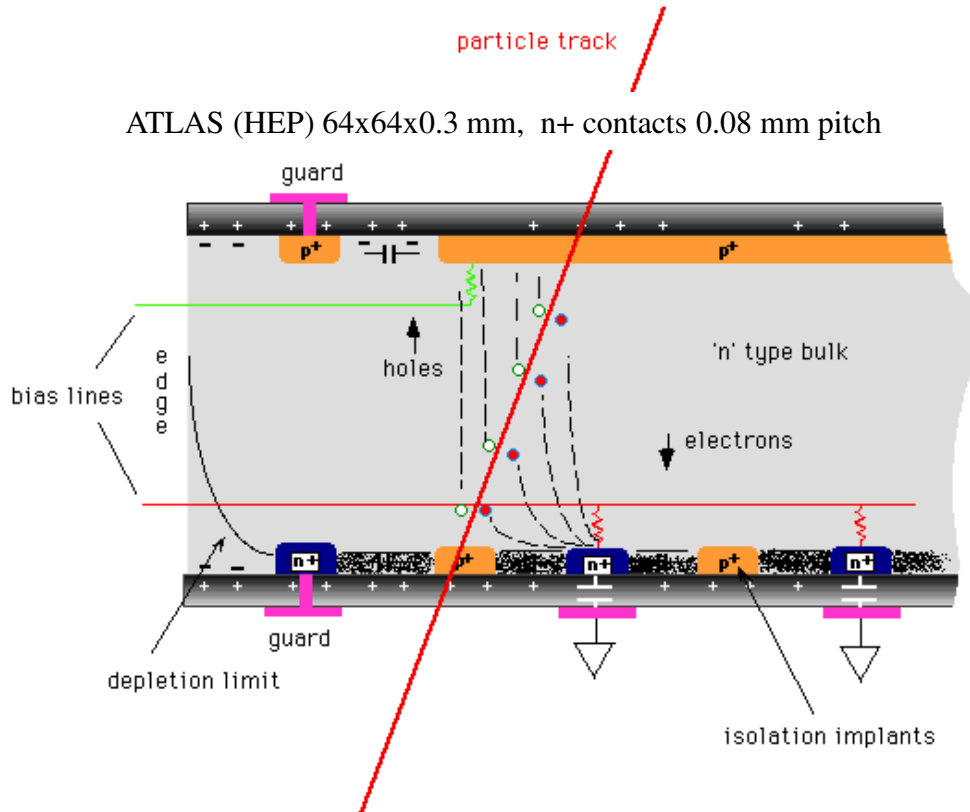


Fig. 11.13 Knoll, 3<sup>rd</sup>, 4<sup>th</sup> Eds.



# Semiconductor Diodes



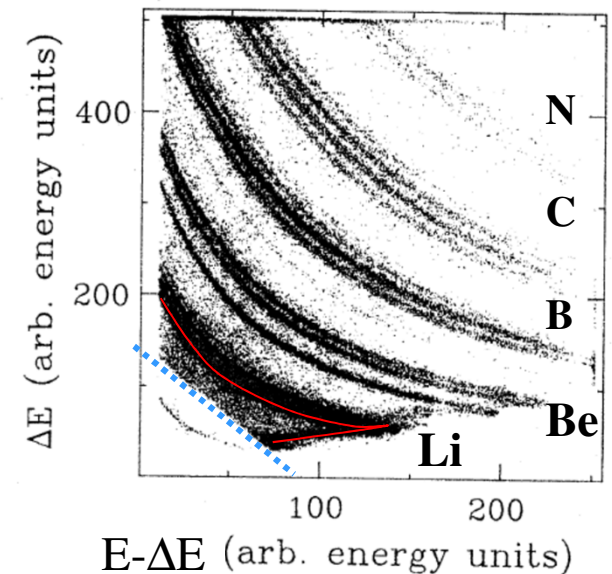
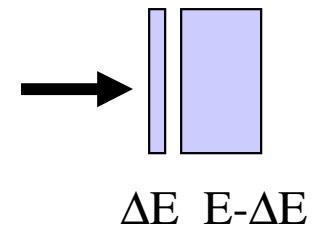
Silicon layers are thin, typically 0.3mm but up 5mm are produced. (dictated by the semiconductor chip industry).

Silicon Detector “telescopes” combine a thin device with a thick device to identify charged particles.

$$\frac{dE}{dx} = C_1 \frac{MZ^2}{E} \ln\left(C_2 \frac{E}{M}\right)$$

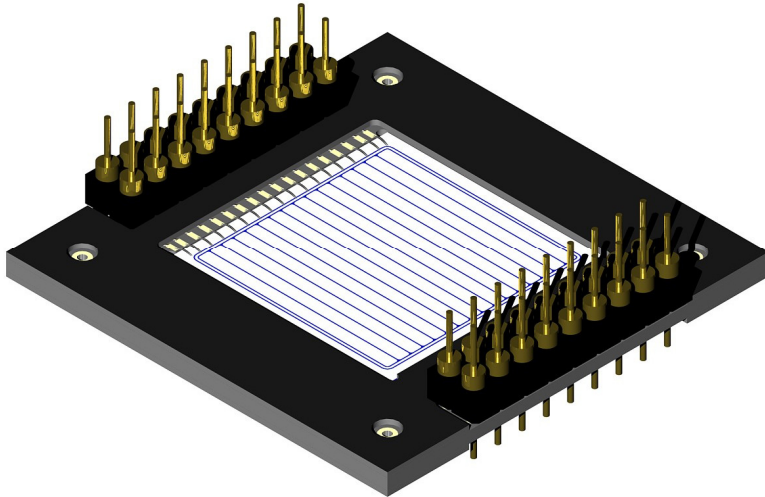
$$\Delta E = \Delta x \left( \frac{dE}{dx} \right) \propto \frac{MZ^2}{E} \rightarrow \Delta E \propto \frac{1}{E}$$

Punch-through  
Software cut



# Semiconductor Diodes – Micron Semiconductor

<http://www.micronsemiconductor.co.uk>



WAFER SIZE	STANDARD SILICON THICKNESSES
	( $\mu\text{m}$ )
4-inch	20, 30, 40, 50, 65, 80, 100, 140, 250, 300, 500, 1000, 1500
6-inch	150, 200, 300, 400, 500, 675, 1000

